

Fig. 3—Phase shift vs applied field. (a) Positively polarized wave. (b) Negatively polarized wave.

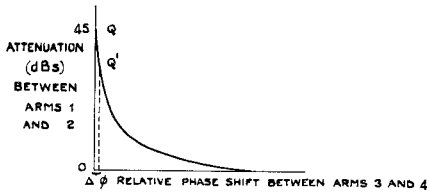


Fig. 4—Attenuation vs relative phase shift.

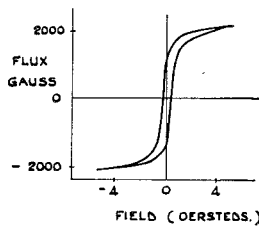


Fig. 5—Hysteresis loop of typical microwave ferrite.

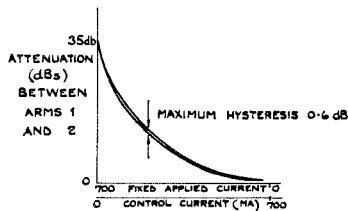


Fig. 6—Attenuator characteristic.

corresponds to the point Q in Fig. 4. If the magnetic field on the ferrite is now changed by a relatively large amount (to point P' in Fig. 3), the actual change in phase in arm 4 is quite small. This results in a small relative phase shift between arms 3 and 4 and this in turn causes the attenuation between arms 1 and 2 to change to the amount corresponding to Q' (Fig. 4). Thus it is obvious that on plotting a curve of attenuation between arms 1 and 2 vs magnetic field applied to the ferrite, a characteristic is obtained whose slope near the maximum attenuation point Q is considerably less steep than that of the curve of Fig. 4. This is shown in Fig. 6.

Since hysteresis is very small near saturation its effect near the steep part of the characteristic of Fig. 6 is very small. Below saturation the hysteresis of the ferrite is more marked (Fig. 5), but since the slope of the characteristic of Fig. 6 is much smaller when the applied field decreases, the effect of this increase in hysteresis is minimized. The final curve for the attenuator is shown in Fig. 6 where it can be seen that the maximum hysteresis measured corresponds to 0.6 db.

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Quarter-Wave Compensation of Resonant Discontinuities*

INTRODUCTION

In designing RF transmission line components it is often necessary to place a short-circuited quarter-wavelength stub in parallel with the line or an open-circuited quarter-wavelength stub in series with the line. The stub can be broadbanded by merely changing the characteristic impedance of the line on either side of the stub for a distance of one quarter wavelength.

BROAD-BAND STUB

It is not generally recognized how broadband a simple stub with quarter-wave transformers can be made. Previous investigators¹ have merely adjusted the transformer impedance for perfect match at two frequencies which depart somewhat from the resonant frequency of the stub without regard for the reflection in the pass band. The following analysis tries to correlate the bandwidth with the allowable reflection in the pass band.

A coaxial broad-band stub is shown in Fig. 1. On each side of the stub the center conductor is enlarged for a length of $\lambda_0/4$ at the center frequency. In these quarter-wave transformers the characteristic impedance is Z_1 . The stub is $\lambda_0/4$ long and its characteristic impedance is Z_2 . The characteristic impedance of the line is taken as 1 ohm in what follows so that Z_1 and Z_2 are multiples of the characteristic impedance.

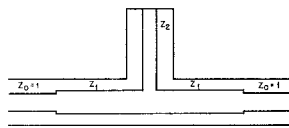


Fig. 1—Coaxial broad-band stub.

The $ABCD$ matrix of the stub plus transformers is

$$\begin{bmatrix} \cos \theta & jZ_1 \sin \theta \\ j(1/Z_1) \sin \theta & \cos \theta \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -j(1/Z_2) \cot \theta & 1 \end{bmatrix} \begin{bmatrix} \cos \theta & jZ_1 \sin \theta \\ j(1/Z_1) \sin \theta & \cos \theta \end{bmatrix}$$

where θ is electrical length of each quarter-wave transformer and the stub. If we let

$\theta = \pi/2 + \phi$, the over-all matrix becomes

$$\begin{bmatrix} -\sin \phi & jZ_1 \cos \phi \\ j(1/Z_1) \cos \phi & -\sin \phi \end{bmatrix} \begin{bmatrix} 1 & 0 \\ j(1/Z_2) \tan \phi & 1 \end{bmatrix} \begin{bmatrix} -\sin \phi & jZ_1 \cos \phi \\ j(1/Z_1) \cos \phi & -\sin \phi \end{bmatrix}$$

which when multiplied gives

$$\begin{bmatrix} \sin^2 \phi - \cos^2 \phi + (Z_1/Z_2) \sin^2 \phi & -jZ_1 \sin \phi \cos \phi (2 + Z_1/Z_2) \\ -j(1/Z_1) \sin \phi \cos \phi (2 - Z_1/Z_2 \tan^2 \phi) & \sin^2 \phi - \cos^2 \phi + (Z_1/Z_2) \sin^2 \phi \end{bmatrix}$$

The insertion loss is given by²

$$\begin{aligned} L &= 10 \log_{10} \{1 + 1/4[(A-D)^2 - (B-C)^2]\} \\ &= 10 \log_{10} \{1 + 1/4[(2/Z_1 - 2Z_1 - Z_1^2/Z_2 \\ &\quad + 1/Z_2) \sin \phi \cos \phi - 1/Z_2 \tan^2 \phi]^2\} \\ &= 10 \log_{10} (1 + m^2/4) \end{aligned} \quad (1)$$

where

$$mZ_2 = R \sin \phi \cos \phi - \tan \phi \quad (2)$$

$$R = 2Z_2/Z_1 - 2Z_1Z_2 - Z_1^2 + 1. \quad (3)$$

A graph of the magnitude of $|m|Z_2$ is shown in Fig. 2.

$R=1$ gives the maximally flat case with a zero derivative at the origin.

For R greater than 1, a triple peaked response is obtained.

Using some simple trigonometric substitutions it can be shown that $\phi_2 = 2\phi_1$; also,

$$m_1Z_2 = \tan \phi_2 \left(\frac{1 - \cos \phi_2}{1 + \cos \phi_2} \right) \quad (4)$$

and

$$R = \frac{2}{\cos^2 \phi_2 + \cos \phi_2} \quad (5)$$

where ϕ_1 is the value of ϕ for worst reflections in the pass band, ϕ_2 is the band edge, and m_1 is the worst value of m in the pass band.

The quantity m_1 is related to the worst voltage standing wave ratio S by

$$m_1 = \frac{S-1}{\sqrt{S}} \quad (6)$$

and the bandwidth is given by

$$BW = 2\phi_2/90. \quad (7)$$

A graph of m_1Z_2 as a function of bandwidth is shown in Fig. 3.

As an example, suppose it is desired to design a stub support for a coaxial line to have a standing wave ratio of no greater than 1.05 over as wide a frequency band as possible. Because of voltage breakdown considerations it is decided that the largest value Z_2 may have is one. Then from (6), $m_1 = 0.0488$, and from Fig. 3 the bandwidth is 70.4 per cent or a frequency ratio of 2.09:1. R is determined from (5), and Z_1 from (3).

The required value of Z_1 for various values of Z_2 is plotted as a function of bandwidth in Fig. 4. This graph shows that the diameter of the quarter-wave transformers is rather critical. The desired Z_1 is only slightly smaller than the zero bandwidth case.

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¹ G. L. Ragan, "Microwave Transmission Circuits," M.I.T. Rad. Lab. Ser., McGraw-Hill Book Co., Inc., New York, N. Y., vol. 9, pp. 173-176; 1948.

² R. M. Fano and A. W. Lawson, "Microwave Transmission Circuits," M.I.T. Rad. Lab. Ser., McGraw-Hill Book Co., Inc., New York, N. Y., vol. 9, ch. 9 and 10; 1948.

